

— FOR IONIZED GAS —

- The heating mainly from free electrons produced by the photoionization of Hydrogen, which is identical in both cases. The free e^- then are able to collisionally excite heavy metals (for which ionization potentials are much higher) which then radiatively de-excite sometimes. Since radiative excitation cross-sections are low, this radiation can escape. This energy ~~is~~ came originally from the colliding free e^- , hence the gas temperature is lowered (cooling).
- Cloud with $10^{10} M_{\odot} H^{\circ}$ (ionization fraction 10^{-4})
 $n_{e^-} = 1 \text{ cm}^{-3}$ neutral gas.

if heating only from kinetic energy from Type II SNe (10^{51} erg per SNe) which occur at a rate of 1 per 100 yr

What is the equilibrium T_{gas} ?

with $\zeta = 10^{-4}$, assume $n = 1 \text{ cm}^{-3}$ refers to H° , also assume negligible metal content for this number.

Use Spitzer Fig 6.2 cooling rates

- At equilibrium, heating rate = cooling rate so need to get rid of

$$\frac{10^{51} \text{ erg}}{10^2 \text{ yr} \times 3.17 \times 10^7 \text{ s/yr}} = \boxed{3.15 \times 10^{44} \text{ erg s}^{-1}} \text{ over entire cloud.}$$

$$\text{Volume of } H^{\circ} = \frac{10^{10} M_{\odot} \cdot 2 \times 10^{33} \text{ g}/M_{\odot}}{1 \text{ cm}^{-3} \times 1.67 \times 10^{-24} \text{ g}/H^{\circ}} = \boxed{1.2 \times 10^{67} \text{ cm}^3}$$

$$\text{Then need to cool at rate of} = \frac{3.15 \times 10^{44} \text{ erg s}^{-1}}{1.2 \times 10^{67} \text{ cm}^3} \times \frac{1}{(1 \text{ cm}^{-3})^2} = \boxed{2.63 \times 10^{-26} \frac{\text{erg}}{\text{cm}^3 \text{ s}}}$$

$$\text{gas Temp from Fig 6.2: } \boxed{T \sim 10^{3.5} = 3 \times 10^3 \text{ K}}$$

E (a) $\frac{\partial u_i}{\partial x_k} = 0$ EXCEPT FOR $\frac{\partial u_x}{\partial x_y}$

THUS

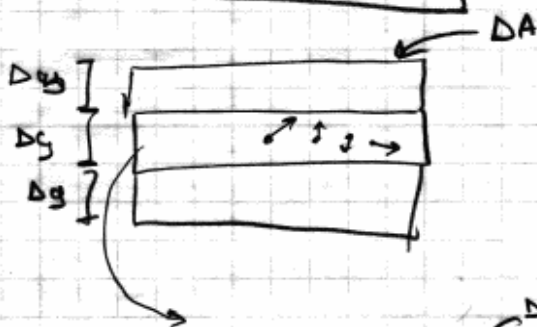
$$D_{xy} = D_{yx} = \frac{\partial u_x}{\partial x_y}$$

ALL OTHER $D_{ij} = 0$

(b) $\mu \vec{\nabla} \cdot \vec{D} = \mu \frac{\partial}{\partial x_k} D_{ij} \Rightarrow$ VISCOUS FORCE $= \mu \rho \vec{\omega}$

$$\text{VISCOUS FORCE} = \mu \frac{\partial u_x}{\partial x_y}$$

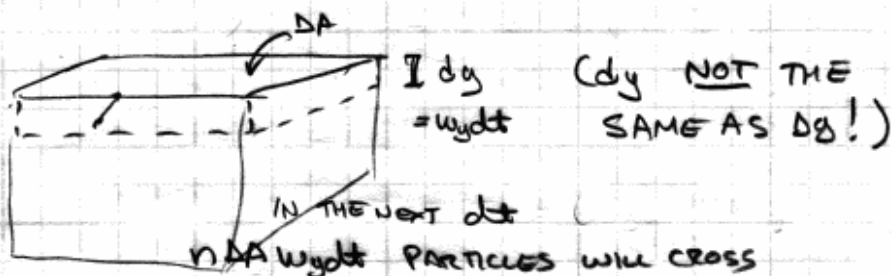
(c)



$$u_0 + \frac{\partial u_x}{\partial y} \Delta y$$

$$u_0$$

$$u_0 - \frac{\partial u_x}{\partial y} \Delta y$$



$$\frac{1}{2} m \langle u_y^2 \rangle = \frac{1}{2} kT$$

WHAT I MEAN BY " u_y " ABOVE $\rightarrow \sqrt{\langle u_y^2 \rangle} = \sqrt{\frac{kT}{m}}$

RATE OF PARTICLES CROSSING $= n \Delta A \Delta y$

$$\text{RATE} = n \Delta A \sqrt{\frac{kT}{m}}$$

(d) TYPICAL MEAN DISTANCE $= \ell =$ MEAN FREE PATH

$$\ell = \frac{1}{n\sigma}$$



(e) IF WE CHOOSE $\Delta y = l$, THEN THE AVERAGE X-VELOCITY OF PARTICLES FROM THE FLUID ELEMENT BELOW IS:

FOR "BELOW"
$$u_x = u_0 - \frac{\partial u_x}{\partial y} l = u_0 - \frac{\partial u_x}{\partial y} \frac{l}{n\tau}$$

THE VELOCITY RELATIVE TO OURS IS

$$-\frac{\partial u_x}{\partial y} \frac{l}{n\tau}$$

SO THE MOMENTUM TRANSFERRED IS

$$-\frac{\partial u_x}{\partial y} \frac{m}{n\tau} \text{ PARTICLES FROM BELOW (EACH ONE)}$$

SIMILARLY

$$+\frac{\partial u_x}{\partial y} \frac{m}{n\tau} \text{ PARTICLES FROM ABOVE (EACH ONE)}$$

(f) TOTAL VISCOUS FORCE ON ONE FACE =

FORCE UNIT AREA $\rightarrow \Delta A \left(\mu \frac{\partial u_x}{\partial y} \right) = (\# \text{ PARTICLES}) \left(\frac{\partial u_x}{\partial y} \frac{m}{n\tau} \right)$

$$\Delta A \mu \frac{\partial u_x}{\partial y} = n \Delta A \sqrt{\frac{kT}{m}} \frac{\partial u_x}{\partial y} \frac{m}{n\tau}$$

$$\boxed{\mu = \frac{\sqrt{mkT}}{\tau}} \quad (\text{TO ORDER OF MAGNITUDE})$$

(g) NOTE: DENSITY DOES NOT APPEAR IN VISCOUS COEFFICIENT! THIS IS REASONABLE SINCE FOR LOWER DENSITY, A PARTICLE CAN CARRY ITS X-MOMENTUM THROUGH A GREATER VELOCITY GRADIENT, SO MOMENTUM TRANSFER IS MORE PER EACH PARTICLE. IN HIGH DENSITY, THEY DON'T GO AS FAR BUT THE COLLISION RATE IS HIGHER, MAKING MOMENTUM TRANSFER FASTER.

